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Reply to the Comment by Martin Rohloff et al.

MICHAEL WILKINSON

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Reply

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MICHAEL WILKINSON

*Department of Mathematics and Statistics, The Open University - Walton Hall,
Milton Keynes, MK7 6AA, England, UK*

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Introduction. – Reference [1] discusses a theory describing experiments on a binary mixture of liquids close to a miscibility transition, with critical temperature T_c . In these experiments periods of turbidity are observed as the system is driven away from the critical point by a slow change of the temperature T . The turbidity is caused by precipitation of droplets which drift towards the interface between the two phases due to their difference in density. Reference [1] presents a theory for the time period Δt between these precipitation events, as a function of the rate of change of the temperature of the system, described by a parameter ξ with dimensions of inverse time. The theory assumes that the initial growth of droplets is by Ostwald ripening (as discussed in [2]), and that this crosses over to collisional growth driven by gravitational settling. The theory predicts that

$$\Delta t \sim K \xi^{-3/7}. \quad (1)$$

The coefficient K is a function of the reduced temperature, $\theta = |T - T_c|/T_c$:

$$K(\theta) = \alpha (D\Lambda\kappa^3)^{-1/7}, \quad (2)$$

where α is a dimensionless constant, D is the interdiffusion coefficient, Λ is a Kelvin length for the phase boundary and the settling speed of a droplet of radius a is $u = \kappa a^2$. The quantities Λ and κ are defined in terms of tabulated material parameters of the mixture, such as the interfacial tension γ and the density difference between the phases, $\Delta\rho$. Because $\Delta\rho$, γ and D all approach zero as a power law in θ at the critical point, it is predicted that $K \sim \theta^{-\eta}$. Using standard critical exponents gives $\eta \approx 0.5$ [1].

In [1] it is asserted that (1) is a good description of the experimental data, for the water/isobutoxyethanol system. This statement is based upon a plot in the PhD

Thesis of Lapp [3], fig. 3.26, which is a scatter plot of $\Delta t \xi^{3/7}$ against temperature. The experimental points collapse onto a curve which shows good agreement with eqs. (1) and (2), over a wide range of values of ξ and θ , with the dimensionless number α close to unity.

The comment by Rohloff, Lapp, and Vollmer, however, is phrased so as to suggest that their experimental results do not support the theory. In the following I argue that the data which they present are fully in accord with the theory as described in [1]. Furthermore, the basis of the claim that there is a disagreement results from an attempt to extend the theory in [1] by arguing that it can be used to determine a bound on the oscillation period. I give two arguments indicating that this proposed extension of the theory is unsound.

Comparison with experimental data. – Figure 1 of the Comment is a version of fig. 3.26 in [3]. I am grateful to the authors of the Comment for making this available in a widely disseminated journal, because it lends experimental support to [1].

In particular, the experimental data in fig. 1 of the Comment indicate that $\Delta t \xi^{3/7}$ is a function of temperature, as the theory predicts, with both ξ and θ varying over approximately two decades. Furthermore, the scatter of data points for different values of ξ appears to be random, indicating that there is no residual systematic dependence upon ξ . This is very strong support for the scaling relation, eq. (1).

Concerning the temperature dependence of $K = \Delta t \xi^{3/7}$, the experimental data points in fig. 1 differ from the theoretical prediction (dotted green curve) by what appears to be a constant offset. This constant offset on a logarithmic scale corresponds to the unknown dimensionless constant α in (2). Moreover, the fact that the data points follow a line of constant slope in

this double-logarithmic plot is an indication that the prefactor K has a power-law dependence on the reduced temperature, $K \sim \theta^{-\eta}$, with $\eta \approx 0.5$. These observations show that $K(\theta)$ is well approximated by eq. (2).

The presentation of the experimental data described in the Comment is a non-trivial test of the theory, and the data in fig. 1 of the Comment are entirely consistent with the prediction in [1].

One aspect of the presentation of fig. 1 of the Comment should be remarked upon. The theoretical curve is not quite the same as that proposed in [1], namely eq. (1) above: an additional factor of $2.4396 \times (4^3/3^4)^{1/7} \approx 2.4$ has been included in the theoretical expression used in the Comment, namely eq. (3) of that work. This exaggerates the apparent “disagreement” with the experimental data, as presented in fig. 1 of the Comment.

Proposed bound on the period. – The model developed in [1] includes a “collision efficiency” ϵ , which was assumed to be of order unity. The authors of the Comment argue that $\epsilon < 1$, and that this inequality leads to a lower bound on the period Δt . They argue that the resulting bound is violated by their experimental data.

There are two flaws in this argument. First, droplet growth in a system which combines gravitational settling and Ostwald ripening is an extremely complicated problem, involving a population of droplets with a distribution of radii $P(a, t)$. This distribution is determined by a supersaturation field satisfying an advection diffusion equation with moving boundaries (due to gravitational settling of the droplets). The simplified model described in eq. (11) of [1], equivalently eq. (1) of the Comment, is a caricature of this complex system in terms of the evolution of the radius of a single droplet, $a(t)$. While this has proven sufficient to surmise scaling properties (as was done in [1]), it is not adequate to establish quantitative bounds on the oscillation period (as proposed in the Comment). Adapting eq. (1) of the Comment to yield a precise bound on the period is an unsound procedure.

Secondly, the argument is based upon the notion that a collision efficiency ϵ cannot exceed unity. This inequality cannot be guaranteed for the collision efficiency which is defined in [1]. The value of ϵ is defined by considering a droplet of radius a falling through a gas of much smaller droplets with volume fraction Φ at a relative speed u . The rate of increase of the volume v_d of the droplet is defined to be

$$\dot{v}_d = \epsilon \pi a^2 u \Phi. \quad (3)$$

In the case where the large droplet grows by coalescence upon contact, the collision efficiency cannot exceed unity

(and it may be substantially smaller, as is thought to be the case for water droplets in clouds [4]). In [1], however, droplets grow by Ostwald ripening, which is effected by diffusive transfer of material between droplets which need not be in contact. A large droplet creates a reduction of the supersaturation field in its neighbourhood, which causes material to evaporate from smaller droplets and condense on the larger droplet [2]. This mechanism allows a falling droplet to collect material from a cylindrical region which has a greater area than its geometrical cross-section, implying that it is possible, in principle, for the collision efficiency defined in (3) to exceed unity.

While the collision efficiencies for water droplets in air have been extensively investigated, those of small droplets in the water/isobutoxyethanol system do not appear to have been studied. Moreover, the physics is fundamentally different (as discussed in [1]) because the experiments are performed close to the critical point (down to $\theta \approx 10^{-3}$). There is, however, an argument supporting the assumption that the collision efficiency ϵ is of order unity. If ϵ were very small or very large, it is unlikely that it would be a constant. Instead, ϵ would be expected to have a systematic dependence upon parameters of the model, which would imply that the predicted scaling relations (1) and (2) would be expected to fail. The success of these predictions is indirect evidence that $\epsilon \approx 1$.

Summary. – The experimental data presented in the Comment are fully consistent with the theory in [1]. Figure 1 of the Comment confirms the $\Delta t \propto \xi^{-3/7}$ scaling, and that the temperature dependence of the prefactor is in accord with the theory. The dimensionless prefactor α would be close to unity if a comparison were made with the theory in the form presented in [1].

The claim that the theory in [1] does not agree with the experimental results is a consequence of an inappropriate attempt to extend it. In particular, the model used in [1] is not sufficiently accurate to support quantitative bounds on the period, and the physics of the growth process does not guarantee that $\epsilon = 1$ is an upper bound on the “collision efficiency”.

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